

Quasi-particle properties of Fermi gases from low density to unitary limits

Bridging nuclear *ab-initio* and density functional theories

Antoine Boulet

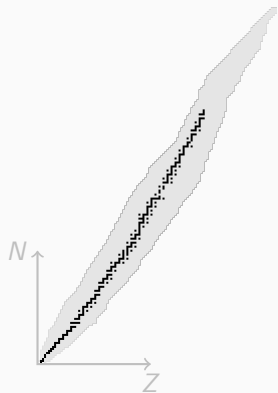
Institut de Physique Nucléaire d'Orsay
antoine.boulet@ipno.in2p3.fr

Collaborator: Denis Lacroix (IPN Orsay)

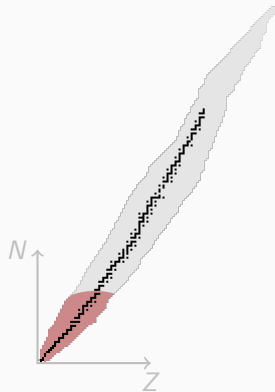
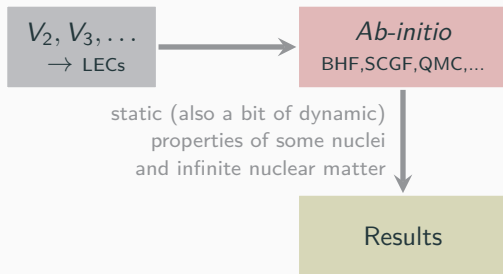
*Symposium on Nuclear Structure and Reactions:
The Next Significant Breakthroughs*
19th march 2019



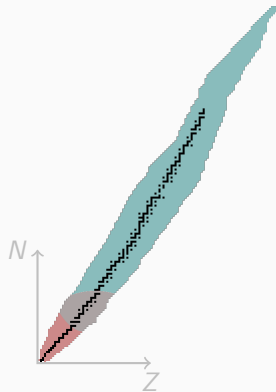
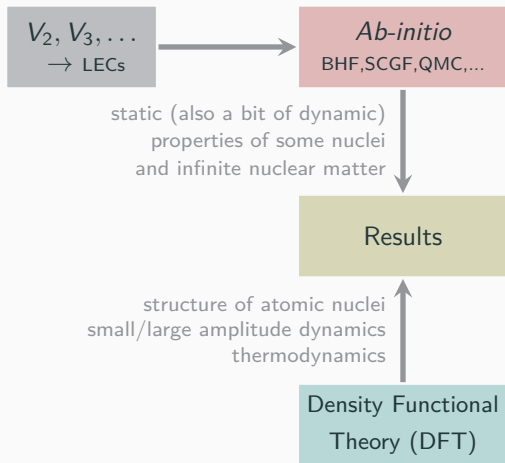
Context and motivation



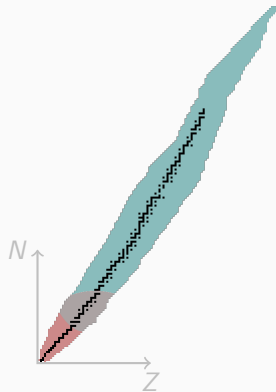
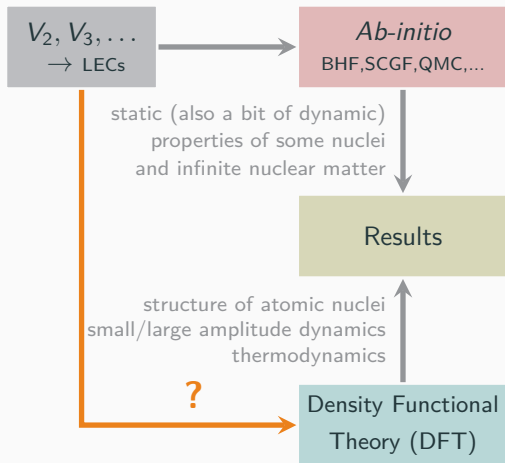
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Context and motivation



**How to relate the bare interaction to DFT
and make it less empirical?**

In this work → a focus on infinite matter

Contents

1. Many-Body Perturbation Theory for dilute Fermi gas in a effective field theory framework
2. Non-perturbative approach: resummation technique

Goal

obtain explicit and simple form for the energy (self-energy) as function of:

- *the density*
- *the low energy constants of the interaction*

The low-density Fermi gas limit: EFT guidance

$$\langle \mathbf{k} | V_{EFT} | \mathbf{k}' \rangle = C_0 + \underbrace{\frac{C_2}{2} (\mathbf{k}^2 + \mathbf{k}'^2)}_{s\text{-wave}} + \dots$$

[Steele and Furnstahl, NPA762 (2000)]

[Beane et al., nucl-th/0008064 (2000)]

[Hammer and Furnstahl, NPA678 (2000)]

$$C_0 = \frac{4\pi}{m} a_s \quad C_2 = \frac{2\pi}{m} a_s^2 r_s$$

Neutron Matter

$$a_s = -18.9 \text{ fm}$$

$$r_s = 2.7 \text{ fm}$$

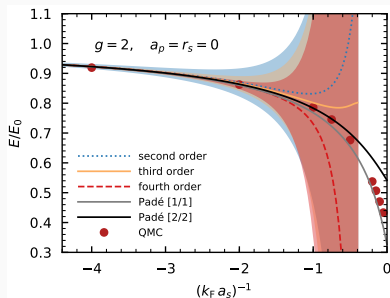
Constructive MBPT

✓ GS energy up to fourth order

[Wellenhofer et al., arXiv (2019)]

UV divergence properly treated

[Kaplan, Savage, Wise, NPB534 (1998)]



Lee-Yang energy density functional

$$\begin{aligned} E(\rho) &= E_{FG} + E^{(1)} + E^{(2)} + \dots \quad \left[E_{FG} = \frac{3}{5} \frac{k_F^2}{2m} \rho \mid \rho = \frac{k_F^3}{3\pi^2} \right] \\ &= E_{FG} \left[1 + \frac{10}{9\pi} (a_s k_F) + \frac{4}{21\pi^2} (11 - 2 \ln 2) (a_s k_F)^2 + \dots \right] \\ &= \frac{3(3\pi^2)^{2/3}}{10m} \rho^{5/3} + \frac{\pi a_s}{m} \rho^2 + \frac{6(11 - 2 \ln 2) a_s^2}{35(3\pi^2)^{-1/3} m} \rho^{7/3} + \dots \end{aligned}$$

✓ analytical dependence in term of ρ and a_s

Difficulties of the perturbative approach

- Perturbative approach valid if $|a_s k_F| \ll 1$
Neutron matter: $a_s = -18.9 \text{ fm} \rightarrow \rho \lesssim 10^{-6} \text{ fm}^{-3} \ll \rho_0 \simeq 0.16 \text{ fm}^{-3}$
- Non perturbative approaches
 - Standard MB techniques: BHF, SCGF, QMC, AFDMC, ...
 - ✓ very powerful
 - ✗ not explicit in $a_s k_F$
 - Resummation technique
 - ✓ analytical in $a_s k_F$ (compatible with a DFT point of view)

Lee-Yang energy density functional

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Basics of diagrammatic framework at zero temperature

[Hammer and Furnstahl, NPA678 (2000)]

$$G(\omega, \mathbf{k}) = \text{---}\text{---}\text{---} = \frac{n_{\mathbf{k}}}{\omega - e_{\mathbf{k}} + i0^-} + \frac{1 - n_{\mathbf{k}}}{\omega - e_{\mathbf{k}} + i0^+}$$

$$\langle \mathbf{k} | V_{EFT} | \mathbf{k}' \rangle = \text{---}\text{---}\text{---} = C_0 \quad [n_{\mathbf{k}} = \Theta(k_F - k) \mid e_{\mathbf{k}} = k^2/2m]$$

Contributing energy diagrams

$$E_{(1)} = \text{---}\text{---}\text{---} \rightarrow (a_s k_F) \rightarrow \textit{Hartree - Fock}$$

$$E_{(2)} = \text{---}\text{---}\text{---} \rightarrow (a_s k_F)^2 \rightarrow \textit{Lee - Yang}$$

$$E_{(3)} = \text{---}\text{---}\text{---} + \text{---}\text{---}\text{---}$$

$$E_{(4)} = \text{---}\text{---}\text{---} + \text{---}\text{---}\text{---} + \text{---}\text{---}\text{---} + \text{---}\text{---}\text{---} + \text{---}\text{---}\text{---}$$

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$$\langle \mathbf{k} | V_{EFT} | \mathbf{k}' \rangle = \text{---} \times \text{---} = C_0 \quad [n_{\mathbf{k}} = \Theta(k_F - k) \mid e_{\mathbf{k}} = k^2/2m]$$

Contributing energy diagrams [Ladder approximation]

$$E_{(1)} = \text{---} \rightarrow \text{---} \rightarrow (a_s k_F) \rightarrow \textit{Hartree - Fock}$$

$$E_{(2)} = \text{---} \rightarrow \text{---} \rightarrow (a_s k_F)^2 \rightarrow \textit{Lee - Yang}$$

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$$E_{(4)} = \text{---} \rightarrow \text{---} + \text{---} \rightarrow \text{---} + \text{---} \rightarrow \text{---} + \text{---} \rightarrow \text{---}$$

[Kaiser, NPA860 (2011)]

Ladder approximation for the energy

Energy resummation

$$E_{int} = \sum_{n=1}^{\infty} \langle \text{diagram} \rangle = \frac{80E_{FG}}{\pi k_F^5} \underbrace{\int_0^{k_F} s^2 ds \int_0^{\sqrt{k_F^2 - s^2}} t dt}_{\text{accessible phase space}} \operatorname{atan} \frac{(a_s k_F) \pi I(s, t)}{\pi - (a_s k_F) R(s, t)}$$

$$E_{int}^{PP} = \sum_{n=1}^{\infty} \langle \text{diagram} \rangle = \frac{80E_{FG}}{\pi k_F^5} \underbrace{\int_0^{k_F} s^2 ds \int_0^{\sqrt{k_F^2 - s^2}} t dt}_{\text{accessible phase space}} \frac{(a_s k_F) \pi I(s, t)}{\pi - (a_s k_F) F(s, t)}$$

[Kaiser, NPA860 (2011)] (no pairing, no self-consistency)

$$F(s, t) = 1 + \frac{s}{k_F} - \frac{t}{k_F} \ln \left| \frac{k_F + s + t}{k_F + s - t} \right| + \frac{k_F^2 - s^2 - t^2}{2sk_F} \ln \left| \frac{(k_F + s)^2 - t^2}{k_F - s^2 - t^2} \right|$$

$$R(s, t) = F(s, t) + F(-s, t)$$

$$I(s, t) = \begin{cases} t/k_F & \text{for } 0 \leq t < k_F - s \\ (k_F^2 - s^2 - t^2)/2sk_F & \text{for } k_F - s \leq t < \sqrt{k_F^2 - s^2} \end{cases}$$

Ladder approximation for the energy

Energy resummation

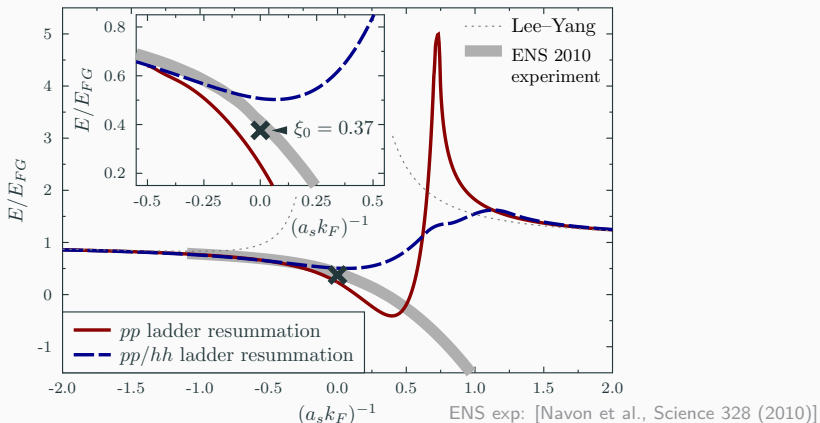
$$E_{int} = \sum_{n=1}^{\infty} \text{diagram} = \frac{80E_{FG}}{\pi k_F^5} \underbrace{\int_0^{k_F} s^2 ds \int_0^{\sqrt{k_F^2 - s^2}} t dt}_{\text{accessible phase space}} \operatorname{atan} \frac{(a_s k_F) \pi I(s, t)}{\pi - (a_s k_F) R(s, t)}$$

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[Kaiser, NPA860 (2011)] (no pairing, no self-consistency)

- ✓ Contains terms to all order in $(a_s k_F)$ in a compact form
- ✓ Expansion in $(a_s k_F) \rightarrow$ Lee–Yang formula
- ✓ Finite limit at unitarity ($a_s \rightarrow \infty$)
- ✗ Implicit function of $\rho = k_F^3/3\pi^2$ (goal: explicit function)

Ladder approximation for the energy



- ✓ correct limit at $a_s k_F \ll 1$ (Lee-Yang expansion)
- ✓ finite limit at unitarity
- ✗ strong dependence of retained diagrams
- ✗ complicated function of $(a_s k_F)$

Phase-Space average Approximation



$$\frac{E_{pp}}{E_{FG}} = 1 + \frac{80}{\pi k_F^5} \underbrace{\int s^2 ds \int t dt}_{\text{accessible phase space}} \frac{(a_s k_F) \pi I(s, t)}{1 - (a_s k_F / \pi) F(s, t)} \xrightarrow{a_s k_F \rightarrow \infty} 0.24$$

PSA of pp ladder resummation = GPS functional

$$\frac{E}{E_{FG}} = 1 + \frac{10}{9\pi} \frac{(a_s k_F)}{1 - (a_s k_F / \pi) \langle F \rangle} \xrightarrow{a_s k_F \rightarrow \infty} 0.32$$

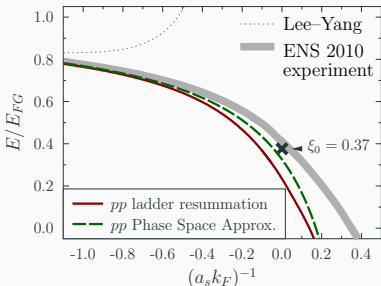
[Heiselberg, PRA63 (2001)] [Schäfer et al., NPA762 (2005)] [Hausmann et al., PRA75 (2007)]

✓ Lee-Yang formula

$$\langle F \rangle = \frac{6}{35} (11 - 2 \ln 2)$$

~ More predictive near unitarity:

$\xi_0 = 0.37$ (accepted value)



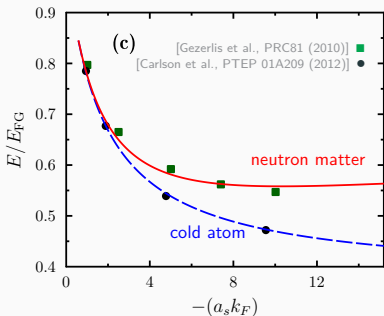
$$\frac{E}{E_{FG}} = 1 + \frac{\frac{10}{9\pi}(a_s k_F)}{1 - \frac{6}{35\pi}(11 - 2 \ln 2)(a_s k_F)} \rightarrow \frac{E}{A} = K + \frac{B\rho}{1 - R\rho^{1/3} + C\rho^{2/3}} + D\rho^{5/3} + F\rho^{\alpha+1}$$

$$\rightarrow 1 + \frac{\frac{10}{9\pi}(a_s k_F)}{1 - \frac{10}{9\pi}(1 - \xi_0)^{-1}(a_s k_F)}$$

[Lacroix, PRA94 (2016)]

[Lacroix, AB, et al., PRC 95 (2017)]

- ✓ unitary limit reproduced
- ✗ Lee-Yang formula

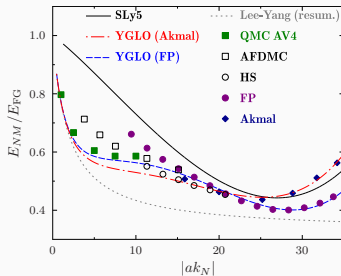


YGLO functional

B, R : Lee-Yang (low density)
→ non-empirical

C, D, F : higher correlations (fit)
→ empirical

[Yang, Grasso, Lacroix, PRC94 (2016)]



Phase-Space average Approximation

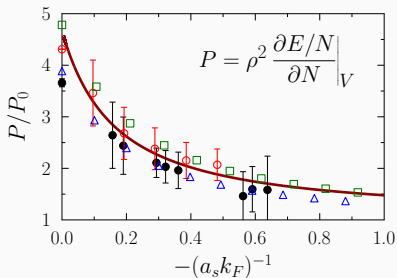


$$\frac{E}{E_{FG}} = 1 + \frac{80}{\pi k_F^5} \underbrace{\int s^2 ds \int t dt}_{\text{accessible phase space}} \operatorname{atan} \frac{(a_s k_F) I(s, t)}{1 - (a_s k_F / \pi) R(s, t)} \Big|_{a_s k_F \rightarrow \infty} = 0.51$$

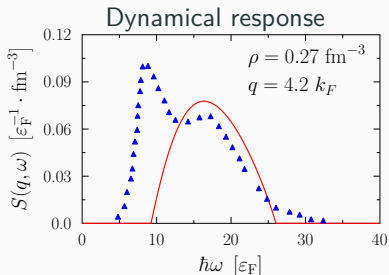
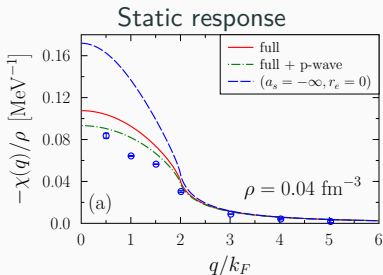
PSA of full ladder resummation = APS functional

$$\frac{E}{E_{FG}} = 1 + \frac{16}{3\pi} \operatorname{atan} \frac{5/24(a_s k_F)}{1 - (a_s k_F / \pi) \langle R \rangle} \Big|_{a_s k_F \rightarrow \infty} = 0.36$$

- ✓ Unitary limit well reproduced
(accepted value: $\xi_0 = 0.37$)
- ✓ Exact Lee–Yang expansion
- ✓ **No adjustment !**



Discussion



✗ effective mass m^*

✗ pairing gap Δ

[AB, Lacroix, PRC97 (2018)]

Goal: extend to self-energy \rightarrow quasi-particle properties (focus on m^*)

Quasi-particle properties: Self-Energy Resummation

Ladder Resummation +
Phase-Space average Approximation

Link with Landau theory of Fermi liquid

$$E_{int} = \sum_{kk'} V_{eff}(k, k') n_k n_{k'}$$

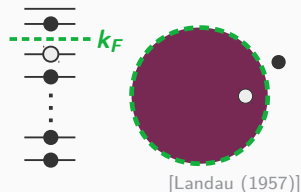
Low-lying
excited states $\downarrow n_k \rightarrow n_k + \delta n_k$

$$\delta E = \sum_k \Sigma^*(k) \delta n_k \mapsto$$

$$\Sigma^*(k) = U(k) + iW(k) = \frac{\delta E}{\delta n_k}$$

Close to
Fermi surface $\downarrow v_{k_F} \equiv \partial_k \epsilon_k|_{k=k_F} \equiv \frac{k_F}{m^*}$

$$\epsilon_k = \epsilon_{k_F} + (k - k_F) \frac{k_F}{m^*} + \dots$$



$$\left[\begin{array}{l} \epsilon_k = \frac{k^2}{2m} + U(k) \\ \frac{1}{2\gamma_k} = -W(k) \end{array} \right]$$

Link with Landau theory of Fermi liquid

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Low-lying
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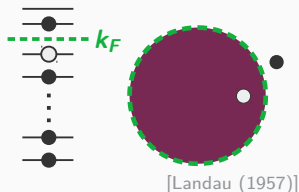
$$n_k \rightarrow n_k + \delta n_k$$

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Close to
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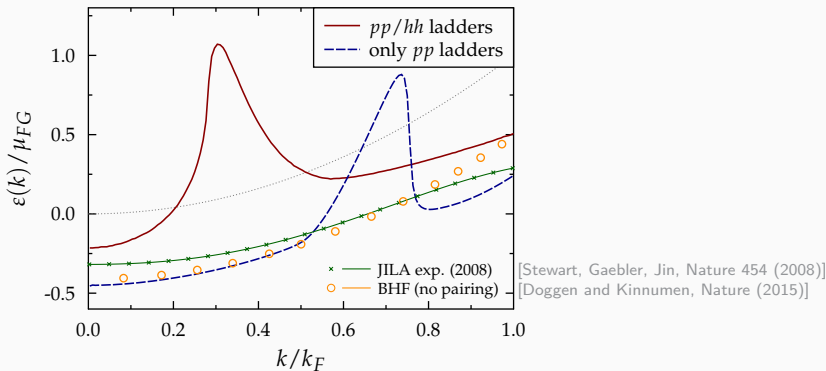


$$\left[\begin{array}{l} \epsilon_k = \frac{k^2}{2m} + U(k) \\ \frac{1}{2\gamma_k} = -W(k) \end{array} \right]$$

Hugenholtz – van Hove theorem (HvH theorem)

$$\mu = E(N + 1) - E(N) = \frac{\partial E}{\partial N} = \epsilon_{k_F}$$

[Hugenholtz, Van Hove, Physica XXIV (1958)]



- ✓ valid at low density \rightarrow Galitskii formula [Galitskii, JETP34 (1958)]:

$$\frac{\varepsilon(k)}{\mu_{FG}} = \frac{4}{3\pi} (a_s k_F) + \phi_2(k) (a_s k_F)^2 + \dots$$

- ✓ finite limit at unitarity ($a_s k_F \rightarrow \infty$)
- ✗ bad prediction for $a_s k_F \gg 1$: pathological
- ✗ strong dependence of retained diagrams

$$E = E_{FG} + \int_{st} \mathcal{E}(s, t)$$

$$\epsilon(k) = \frac{k^2}{2m} + \int_{st} \mathcal{U}(s, t, k)$$

Strategy

$$E = E_{FG} + \int_{st} \mathcal{E}(s, t)$$



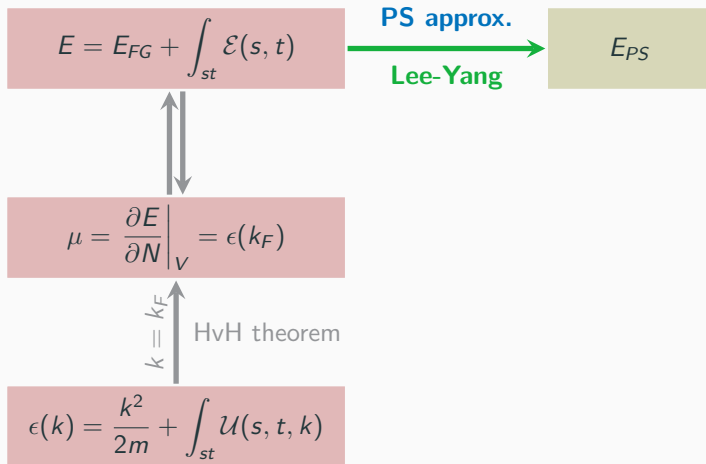
$$\mu = \left. \frac{\partial E}{\partial N} \right|_V = \epsilon(k_F)$$



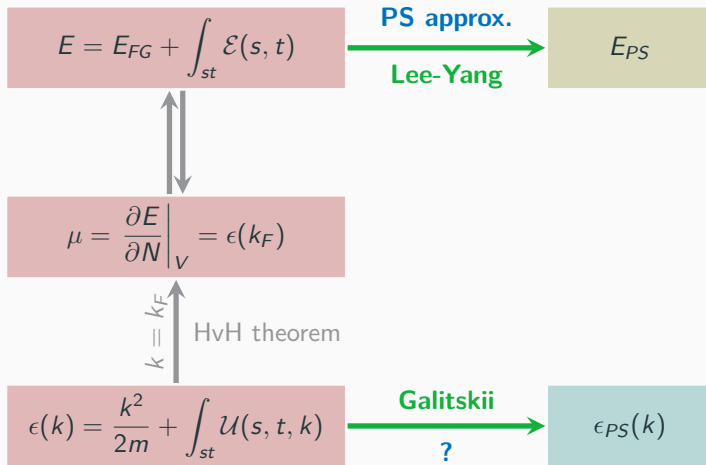
HvH theorem

$$\epsilon(k) = \frac{k^2}{2m} + \int_{st} \mathcal{U}(s, t, k)$$

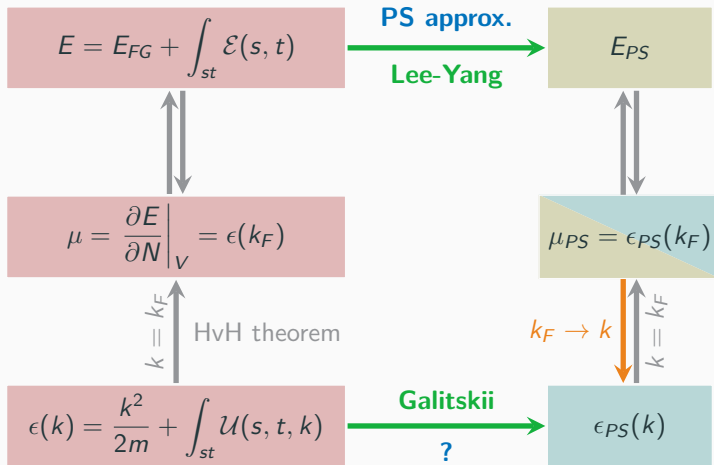
Strategy



Strategy



Strategy




Phase-space average approximation: GPS case

$$\frac{E}{E_{FG}} = 1 + \frac{10}{9\pi} \frac{(a_s k_F)}{1 - (a_s k_F / \pi) \frac{9\pi^2}{14} \phi_2(k_F)}$$

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$$\mu = \left. \frac{\partial E}{\partial N} \right|_V$$



$$\frac{\mu}{\mu_{FG}} = 1 + \frac{4}{3} \frac{(a_s k_F)}{\pi - (a_s k_F) \frac{9\pi^2}{14} \phi_2(k_F)} + \frac{2}{9} \frac{(a_s k_F)^2 \frac{9\pi^2}{14} \phi_2(k_F)}{[\pi - (a_s k_F) \frac{9\pi^2}{14} \phi_2(k_F)]^2}$$

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$$\phi_2(k_F) \rightarrow \phi_2(k)$$


$$\frac{\epsilon(k)}{\mu_{FG}} = 1 + \frac{4}{3} \frac{(a_s k_F)}{\pi - (a_s k_F) \frac{9\pi^2}{14} \phi_2(k)} + \frac{2}{9} \frac{(a_s k_F)^2 \frac{9\pi^2}{14} \phi_2(k)}{[\pi - (a_s k_F) \frac{9\pi^2}{14} \phi_2(k)]^2}$$

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$$\mu = \left. \frac{\partial E}{\partial N} \right|_V \quad \checkmark \text{ Lee-Yang Formula}$$

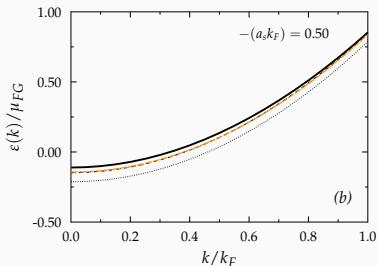
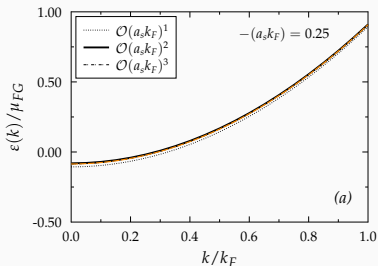
$$\frac{\mu}{\mu_{FG}} = 1 + \frac{4}{3} \frac{(a_s k_F)}{\pi - (a_s k_F) \frac{9\pi^2}{14} \phi_2(k_F)} + \frac{2}{9} \frac{(a_s k_F)^2 \frac{9\pi^2}{14} \phi_2(k_F)}{[\pi - (a_s k_F) \frac{9\pi^2}{14} \phi_2(k_F)]^2}$$

$$\phi_2(k_F) \rightarrow \phi_2(k) \quad \checkmark \text{ HvH theorem } \mu = \epsilon(k)$$

$$\frac{\epsilon(k)}{\mu_{FG}} = 1 + \frac{4}{3} \frac{(a_s k_F)}{\pi - (a_s k_F) \frac{9\pi^2}{14} \phi_2(k)} + \frac{2}{9} \frac{(a_s k_F)^2 \frac{9\pi^2}{14} \phi_2(k)}{[\pi - (a_s k_F) \frac{9\pi^2}{14} \phi_2(k)]^2}$$

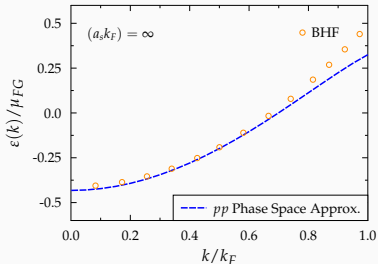
✓ Galitskii Formula

Single particle energy



MBPT: [Platter *et al.*, NPA714 (2003)]
[AB, Lacroix, submitted to J. Phys. G]

- ✓ exact expansion up to $(a_s k_F)^2$
→ Galitskii formula
- ✓ pathologies removed for
 $a_s k_F \gg 1$ (better prediction)
- ✓ simpler function of the density

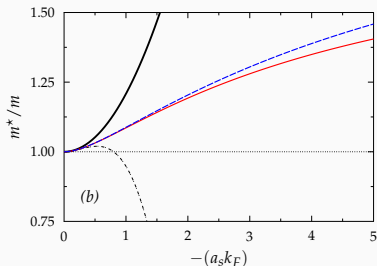
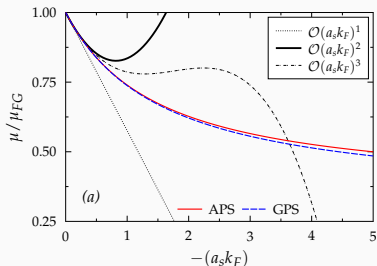


BHF: [Doggen and Kinnunen, Nature (2015)]

Chemical potential and effective mass

$$\mu = \epsilon(k_F)$$

$$\frac{m}{m^*} = \frac{m}{k_F} \left. \frac{\partial \epsilon_k}{\partial k} \right|_{k_F}$$



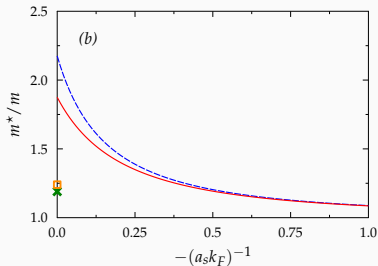
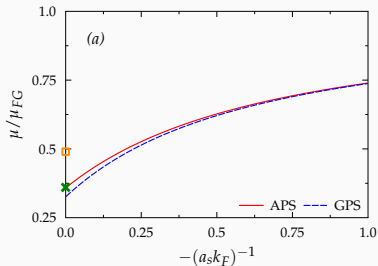
MBPT: [Platter *et al.*, NPA714 (2003)]
[AB, Lacroix, submitted to J. Phys. G]

- ✓ Expansion valid up to $(a_s k_F)^2 \rightarrow$ Galitskii formula
- ✓ Simple and explicit dependence in density
- ✓ Finite limit at Unitarity

Discussion

$$\mu = \epsilon(k_F)$$

$$\frac{m}{m^*} = \left. \frac{m}{k_F} \frac{\partial \epsilon_k}{\partial k} \right|_{k_F}$$



	superfluid \times (expected)	normal \square (expected)	GPS	APS
(at unitarity) μ/μ_{FG}	0.37	0.49	0.32	0.36
m^*/m	1.19	1.24	2.18	1.88

It seems difficult to enforce unitarity without adjustment

Summary

- Ladder (pp or pp/hh) resummation from E to $\Sigma^*(k)$
 - ✗ quite complex density dependence
 - ✗ strong dependence on the selected diagrams
- **Phase-space approximation of the energy**
 - ✓ simple and explicit density dependence
 - ✓ predictive from low density to unitarity without adjustment
- **Phase-space approximation of the self-energy**
 - ✓ simple and explicit density dependence
 - ✓ predictive at low and intermediate density
 - ✗ Unitary limit far from expected results: need to be adjusted
 - ✗ Pairing effect: from normal to superfluid

Perspective

Make explicit the link with density functional theory

→ apply to finite systems